DEMONSTRATION OF TERRASAR-X SCANSAR PERSISTENT SCATTERER INTERFEROMETRY

Fernando Rodriguez Gonzalez (1), Ramon Breic (1), Nestor Yague-Martinez (1,2), Robert Shau (1), Alessandro Parizzi (1), Nico Adam (1)

(1) Remote Sensing Technology Institute (IMF), German Aerospace Center (DLR) Oberpfaffenhofen, 82234 Weßling, Germany. Email: Fernando.RodriguezGonzalez@dlr.de
(2) Chair of Remote Sensing Technology, Technical University Munich (TUM) Arcisstraße 21, 80333 Munich, Germany

ABSTRACT
The concept of PSI Wide Area Product was introduced in ESA’s GMES Terrafirma project [1] in preparation for the upcoming Sentinel-1 mission. Sentinel-1 provides unprecedented capabilities for wide area mapping due to the coverage of the IW (Interferometric Wide swath) mode, which is the predefined mode over land. The extension of the PSI technique to wide area deformation mapping presents many technical challenges, especially in non-urban areas and in mountainous areas [2]. The Integrated Wide Area Processor (IWAP) [3] has been developed at DLR for this purpose. This paper presents the processor extensions performed in order to support the ScanSAR and TOPS SAR acquisition modes in preparation for the Sentinel-1 mission.

1. INTRODUCTION
The Sentinel-1 system has been conceived in order to provide wide swath SAR interferometric capabilities. A wide swath is achieved by means of the novel TOPS (Terrain Observation by Progressive Scan) SAR acquisition mode [4]. In this mode the radar beam cyclically scans a group of neighbouring swaths, as it is also performed in the ScanSAR acquisition mode. Furthermore, in TOPS the antenna is systematically steered forwards during the scan of a burst, eliminating thus the main drawbacks of ScanSAR, namely the effects of scalloping and the azimuth variant noise equivalent $\sigma_n$ and signal-to-ambiguity ratio.

Nevertheless, the spectral characteristics of the focussed SAR complex images (linear Doppler Centroid variation over azimuth) and the interferometric processing techniques are analogous for ScanSAR and TOPS SAR [5][6]. The main difference is the higher azimuth co-registration accuracy required for TOPS SAR, which is due to its higher Doppler Centroid variations within a burst caused by the antenna forward steering.

In the framework of ESA GMES Terrafirma project (ESRIN/Contract no. 4000109669/13/I-AM) as well as in order to support ESA in its Sentinel-1A Commissioning Phase activities (ESA Contract No. 4000111074/14/NL/MP/II) [7], DLR’s IWAP processor has been extended to support both burst acquisition modes: ScanSAR and TOPS SAR.

During the Terrafirma project there were no Sentinel-1 stacks available. As a consequence, the processing demonstration has been performed with a stack of 65 TerraSAR-X ScanSAR images. It covers an area of approximately 100x100 km surrounding Lake Mead, an artificial water reservoir impounded by the Hoover Dam. The scene is characterised by high topography variations, from 200 to 2400 meters. The stack was acquired from August 2008 to January 2011 and its master on the 10th December 2009. During this period, the water elevation of Lake Mead progressively descends from approximately 337.5 m to 332.5 m [8] (see Fig. 1-left). Seasonal water elevation changes can be also clearly seen. The time baseline distribution of the stack is presented in Fig. 1-right.

2. METHODOLOGY

InSAR Processing
The principles of the implemented burst mode interferometric processing have been presented in [6]. The co-registration approach is DEM based geometric co-registration with correction of a residual offset (both in range and in azimuth) per interferogram determined either by Incoherent Cross Correlation (ICC) or by Enhanced Spectral Diversity (ESD) [5]. In this test case the offset has been determined by ICC, while ESD has been used for quality control.

Selected interferograms are depicted in Fig. 2-left and center in order to illustrate the atmospheric conditions within the stack. High turbulence and stratification can be observed. The acquisitions have been performed in the late afternoon (ascending orbit), when atmospheric turbulence is higher. The high atmospheric stratification effect is due to the big height variations within the scene, causing approximately two fringes in the example shown.

PSI Processing
The approach designed for PSI processing is conditioned by the elongated shape of the bursts, which have an azimuth length on the order of only...
2.3 km. Strategies based on burst per burst processing and mosaicking of burst level PSI estimates are not feasible. Our approach is to process all bursts jointly.

In order to perform PSI detection the scene is partitioned into non-overlapping areas, each corresponding to the central part of the bursts. PSs are detected within each of these burst areas according to an SCR threshold and sub-pixel position determination. Arcs are generated from burst to burst, including beam to beam. Thus it is useful to review the impact of the Doppler Centroid in the arc double differences. The SAR impulse response function in zero Doppler geometry in the case of an observation squint angle $\beta$ is given by [9]

$$s(t, r) = C \cdot s_r(t - t_0) \cdot s_a \left( t - t_0 - \frac{c \cdot \tan \beta}{v} \right) \cdot \exp \left( -j \frac{4\pi}{\lambda} r_0 \right) \cdot \exp \left( j2\pi f_{DC} (t - t_0) \right) \cdot \exp \left( -j2\pi f_0 (1 - \cos \beta)(t - t_0) \right)$$

(1)

where $t, r$ are the fast and slow time, $t_0, r_0$ are the closest approach times and distance, $C$ is a complex constant, $s_r$ and $s_a$ the compressed range and azimuth envelopes, $v$ effective platform forward velocity, $f_0$ the radar frequency, $f_{DC}$ the Doppler Centroid and $\lambda$ the wavelength. In the presence of a small squint the last exponential term can be neglected. According to this, the phase of a PS $l$ in the master scene extracted with an azimuth subpixel error $\Delta t_{PTA1}$ is given (up to a constant) by

$$\phi^M_l = -\frac{4\pi}{\lambda} r_0^M + 2\pi f_{DC,1} \Delta t_{PTA1}$$

(2)

and the phase of the same point extracted from a slave image with an azimuth co-registration error of $\Delta t_{CO}$ is given by

$$\phi^S_l = -\frac{4\pi}{\lambda} r_0^S + 2\pi f_{DC,1} (\Delta t_{PTA1} - \Delta t_{CO})$$

(3)

where the super-index $M$ and $S$ denote master and slave images. Let $\phi^M_i = \phi^M_l - \phi^M_j$ be the interferometric phase of a point $i$. Then the phase of a double difference between points $l$ and $j$ is given by

$$\phi^M_i - \phi^M_j = -\frac{4\pi}{\lambda} (\Delta r_{0,1}^M - \Delta r_{0,2}^M) - 2\pi (\Delta f_{DC,1}^M - \Delta f_{DC,2}^M) \Delta t_{PTA1}$$

$$+ 2\pi (f_{DC,1}^S - f_{DC,2}^S) \Delta t_{PTA2} + 2\pi (f_{DC,1}^S - f_{DC,2}^S) \Delta t_{CO}$$

(4)
where the notation $\Delta[\mathcal{M}_{i}^{M-S} = \mathcal{M}_{i}^{M} - \mathcal{M}_{i}^{S}$ denotes the difference between master and slave acquisitions for the point $i$. The first term is the one we exploit for inverting difference of deformation and residual height. The second term is due to the subpixel position determination errors. Nevertheless its proportionality factors $\Delta f_{DC,i}$ are very small for current SAR missions due to accurate burst synchronization and attitude control. By performing some prior oversampling and/or point target analysis this term can be minimized. The third term is much more important, since it depends on the difference of Doppler Centroids between points $1$ and $2$, which for ScanSAR and TOPS SAR can be very high (especially in the case of arcs between different bursts). This term is minimised by means of an accurate co-registration in the stacking process. Thus, as long as the co-registration has been done accurately, it is possible to use any standard PSI estimator.

Due to the high height variation in this test site it is required to mitigate the stratified atmospheric delay in order to i) avoid biased deformation components correlated with height and to ii) generate reliable arcs between different heights for a robust reference network construction and inversion. The approach used here has been iterative: based on an initial PS processing, the unwrapped PS phases have been used to fit a linear height-to-phase model in the mountain located in the central area and its surroundings. Note that this data based approach has the risk of compensating deformation components which are correlated with height. Weather model data based approaches would avoid this risk.

The further PSI processing has been performed using standard techniques described in [2][3].

3. RESULTS

An overview of the deformation estimates is provided in Fig.3-left. A spatial ramp has been compensated, reducing thus low pass residual components. Nevertheless, still residual low pass atmospheric signals can be seen in the results. This may be due to the high atmospheric turbulence as well as to residuals from the stratification mitigation and to the time span of the stack.

Further from these effects, we observe a big subsidence pattern in the northern area of Lake Mead (Fig.3-center). In [10] it was also observed that this area did not follow the global deformation pattern. The explanation that the authors provided [10] was that the local pattern possibly was caused by the poroelastic response of a sedimentary layer with groundwater communicated with the lake. The authors explain that a
high lake level may induce a poroelastic uplift of the surface of the sedimentary layer. Since in our acquisition time span the water level goes down, we observe subsidence, which would be consistent with the explanation in [10].

A subsidence which a priori seems a landslide can be observed in Fig. 3-right. Further from these phenomena, coastal subsidence around the basin can be observed on selected areas, as shown in Fig. 3-left.

The topographic phase compensation has been performed with the SRTM DEM, which was acquired in February 2000. The master scene was acquired in December 2009. During this time, the mean Lake Mead water level went down from approximately 370 to 334 meters [8]. As a consequence, areas that used to be under the water in 2000, were new coastal areas in 2009. This phenomenon can be clearly observed in the topographic update estimates. In Fig.4-left green means no topographic update, whereas red-orange-yellow is negative update, i.e. regions that before used to be under the water, as the small island.

Gaps in the estimation results can be observed in Fig.4-left. These are due to the burst synchronization and to the PSI requirement that we see the target in all or most of the images. If enforced that a point must be visible in all images, then a gap of 125-150 meter would have no estimates. For this visualization this condition was relaxed to PS visible in at least 59 images.

Figure 4-center shows the average PS time coherence calculated on a 100x100 grid. Points in black mean that a minimum of 10 PSs was not reached. The coherence is homogeneously distributed, with some local variations which seem to be correlated with topography. Figure 4-right shows histograms of PS time coherence for 200 Hz Doppler Centroid bins, with the local Doppler Centroid of the bin used as colour code. It can be seen that the distribution of coherence is almost the same independently of the Doppler Centroid.

4. EXPLOITATION OF OVERLAP AREAS
The ScanSAR and TOPS SAR acquisition modes have overlap between bursts within a beam as well as between beams. ESD exploits these overlaps during co-registration in order to accurately estimate the azimuth co-registration offset. As already proposed in [5], it is possible to exploit the ESD technique with the PSs interferometric phases. One advantage of PSs is that they do not decorrelate (or very slowly) in time and thus they can be exploited even with master and slaves images having a high temporal separation, i.e. when the coherence of distributed scatterers decreases due to natural decorrelation phenomena [11].

The approach we suggest is to exploit the PSs at the overlap areas after PSI detection in order to perform a quality control of the azimuth co-registration accuracy. As previously mentioned, PS detection was performed on a partition composed of the central areas of the bursts. We will call these detected PS “primary PS”, since they have been detected due to their SCR and sub-pixel position determination has been performed in their mean amplitude map. For the other areas of the bursts, instead of performing again an independent detection, PSs are extracted in the same positions that PSs were detected in one of the neighbouring bursts. This extra PSs will be called “secondary PSs”, since they have not been “detected” directly themselves. It follows that each pair of primary-secondary PSs should provide two observations at different squints from the same target. This concept is visualized in Fig.5-left. Depending on the burst combinations, we obtain primary-secondary pairs with Doppler Centroid differences around 2000 Hz (intra beam), 1500 Hz (inter beam) and 500 Hz (inter beam). The locations of the two first groups are depicted in Fig.5-center and right. Note that the first beam, which was acquired with a lower range bandwidth, has been excluded for the analysis of the primary-secondary PSs.

Firstly the nature of this pairs has been analysed. An evaluation of the coherence of primary versus secondary PSs for the 1500 Hz group is shown in Fig.6-left. The coherence of the secondary PS is mostly lower than that of the primary. The histogram of coherence-loss versus power-loss for primary-secondary PSs shown in Fig.6-center indicates that this lower coherence of the secondary PS could be explained by a power loss. This would be an indicator that PSs in this dataset have a spatial extent and directivity, which is actually consistent with the fact that there are almost no man-made structures. By observing these PSs with different Doppler Centroids we observe different parts of the backscattering pattern and thus different backscattering power. Since the primary PSs are detected on the basis of their SCRs, it is reasonable that they have on average a higher power and coherence.

The PSs that we are exploiting in this dataset are possibly mostly stable rocks. This is further supported by the increase of the standard deviation of the PSs residual phase (jointly evaluated for each interferogram) as the time separation from master increases, which is shown in Fig. 6-right. The standard deviation has been evaluated with primary PSs with coherences between 0.7 and 0.9. This behaviour is an indicator of decorrelation phenomena, which is reasonable given the short wavelength of X-band.

The behaviour of the three primary-secondary groups is surprisingly similar and requires further analysis. A comparison of the deformation estimates for primary-secondary pairs exhibits biases in the order of hundredths of mm/year deformation and a hundredth of m residual topography.

Finally, a preliminary evaluation of the applicability of these primary-secondary PS pairs for azimuth co-registration quality control is illustrated in Fig.7.
this purpose, an interferogram with high co-registration error has been selected. On the left side, the histogram of the standard ESD residual phase after interferogram generation (distributed scatterer, intra beam) is depicted. The histogram of the differential PS primary-secondary phases for the 2000Hz and 1500 Hz Doppler difference groups are respectively shown in Fig.7 -center and right. The phase difference has been calculated so that the difference of Doppler is consistently positive in order to be a direct indicator of residual co-registration offset.
Initial analysis on the full stack indicates a good consistency between the average differential phases of the three Doppler difference groups and the average differential phase of standard ESD. Nevertheless, for the 2000 Hz group, mean phase mismatches with respect to standard ESD of maximum value 1.3° can be observed. These differences could be partially due to the spatial variation of the azimuth residual shift and the different spatial sampling for standard and PS ESD. For the 1500 Hz and the 500 Hz some higher mismatches are found, which require further investigation and could be problematic due to the higher phase-to-azimuth-shift conversion factors.

As a consequence, our current concept is to exploit mainly the 2000 Hz Doppler Centroid difference group, while the other two groups could be used for the verification of the azimuth co-registration ambiguity band. This approach has already been demonstrated for TOPS SAR co-registration purposes using distributed scatterers [12]. The practical evaluation with PS ESD must still be performed.

5. ACKNOWLEDGEMENTS
The WAP is developed within the ESA project Terrafirma with the ESRIN/Contract no. 4000109669/13/I-AM.

6. CONCLUSION
The extension of DLR’s IWAP to support PSI with ScanSAR and TOPS SAR acquisition modes has been presented.

The extension has been demonstrated with a TerraSAR-X ScanSAR data stack in the area surrounding Lake Mead. Subsidence phenomena and a landslide in Lake Mead basins’ have been identified. The impact of the water level change on the residual topography with respect to SRTM has been illustrated.

A preliminary evaluation on the behaviour of PSs observed from different bursts and their potential for co-registration quality control has been performed. Further research on urban and TOPS SAR cases should be performed.

7. REFERENCES